

EQUILIBRIA IN THE SYSTEM  $2N_2 + H_2$  AND  $N_2 + 8H_2$ ,  
AND THERMODYNAMIC PROPERTIES OF  $N_2H_4$  AT  $1000^\circ$  TO  $5000^\circ K$

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Although high temperature equilibrium data on systems containing carbon, hydrogen, and nitrogen have been reported, the same type of data for the more simple system containing hydrogen and nitrogen<sup>1</sup> have not been reported. Both systems are of obvious interest in the chemistry of reactions in a high-temperature plasma.

We presently consider the equilibria in the systems  $2N_2 + H_2$  and  $N_2 + 8H_2$  at  $1000^\circ$  to  $5000^\circ K$ . Ideal gas behavior is assumed. The following species only are considered to be present in the gas:  $H_2$ ,  $N_2$ ,  $H$ ,  $N$ ,  $NH$ ,  $NH_2$ ,  $NH_3$ ,  $N_2H_4$ . Undoubtedly, there are other species and ions present, such as  $N_2H_3$ , but the concentrations of these will be negligible compared to those considered here. It is expected that  $N_2H_4$  will also be negligible, but we include this species because of its practical interest.

Eight equations are required in order to determine the concentrations of the eight components. Five equations express the equilibrium constants for the formation of  $H$ ,  $N_2$ ,  $NH$ ,  $NH_2$ , and  $NH_3$ , where the data are taken from the JANAF Thermochemical Tables.<sup>2</sup> The sixth equation, the equilibrium constant for formation of  $N_2H_4$ , is calculated herein. The seventh equation expresses the total pressure as the sum of the partial pressures,  $\sum_i P_i = P$ . The eighth equation expresses the material balance; i.e., for the system  $2N_2 + H_2$

$$2 P_{H_2} + P_H + 0.5 P_{NH} + 1.5 P_{NH_2} + 2.5 P_{NH_3} + 3 P_{N_2H_4} \\ - P_{N_2} - 0.5 P_N = 0,$$

and for the system  $N_2 + 8H_2$

$$8 P_N + 16 P_{N_2} + 7 P_{NH} + 6 P_{NH_2} + 5 P_{NH_3} + 12 P_{N_2H_4} \\ - P_H - 2 P_{H_2} = 0.$$

The set of eight equations were solved for the different temperatures and pressures of one and a thousand atmospheres on a computer. A single pass in the computation was sufficient. Figures 1 to 4 summarize the results.

The partial pressure of  $N_2H_4$ , as suspected, is negligible in all of the systems, although it goes through a maximum. The partial pressure of  $NH_3$  decreases rapidly with increasing temperature, showing the low stability of this species at the high temperatures. It is interesting to note the stability of  $N_2$  even at  $5000^\circ C$ , as evidenced by the small decrease in partial pressure of this species from  $1000^\circ$  to  $5000^\circ K$ .

There is only a very small effect on the concentration of the species  $NH$ ,  $NH_2$ ,  $N_2H_4$ , and  $NH_3$  by varying the ratio of  $N_2:H_2$ . However, an increase in pressure from one to 1000 atmospheres does cause a slight shift in equilibrium, so that the relative concentrations of  $NH_3$ ,  $N_2H_4$ , and  $NH_2$  increases slightly.

#### Thermodynamic Properties of $N_2H_4$

Scott, et al,<sup>3</sup> have reported thermodynamic properties of  $N_2H_4$  to  $1500^\circ K$ . We extend these data to  $5000^\circ K$  by using the more recent spectroscopic data of Giguere and Liu.<sup>4</sup> The latter workers were unable to determine the symmetrical stretching ( $\nu_2$ ) and torsional ( $\nu_7$ ) vibrations, so that we use the values of Scott, et al, for these two frequencies. The three components of the moment of inertia, the potential barrier to internal rotation, the reduced moment calculated by Pitzer's formula,<sup>5</sup> and the heat of formation at absolute zero are those given by Scott. Table I summarizes the spectroscopic data.

In the calculation of the equilibrium constants of  $N_2H_4$ , the free energy functions of  $H_2$  and  $N_2$  were obtained from the tables of the U.S. Bureau of Standards.<sup>6</sup> Table II summarizes the thermodynamic data.

The values of the thermodynamic functions calculated by us differ from those of Scott<sup>3</sup> by 1.5 to 3.5% in the different functions in the range  $1000^\circ$  to  $1500^\circ K$ , since the vibrational contributions in this work are based on the more recent data of Giguere and Liu.<sup>4</sup> The free energy of formation and  $-\log K_p$  for this temperature region are about 1.5% and 2% higher, respectively, than those of Scott and co-workers.

#### Acknowledgement

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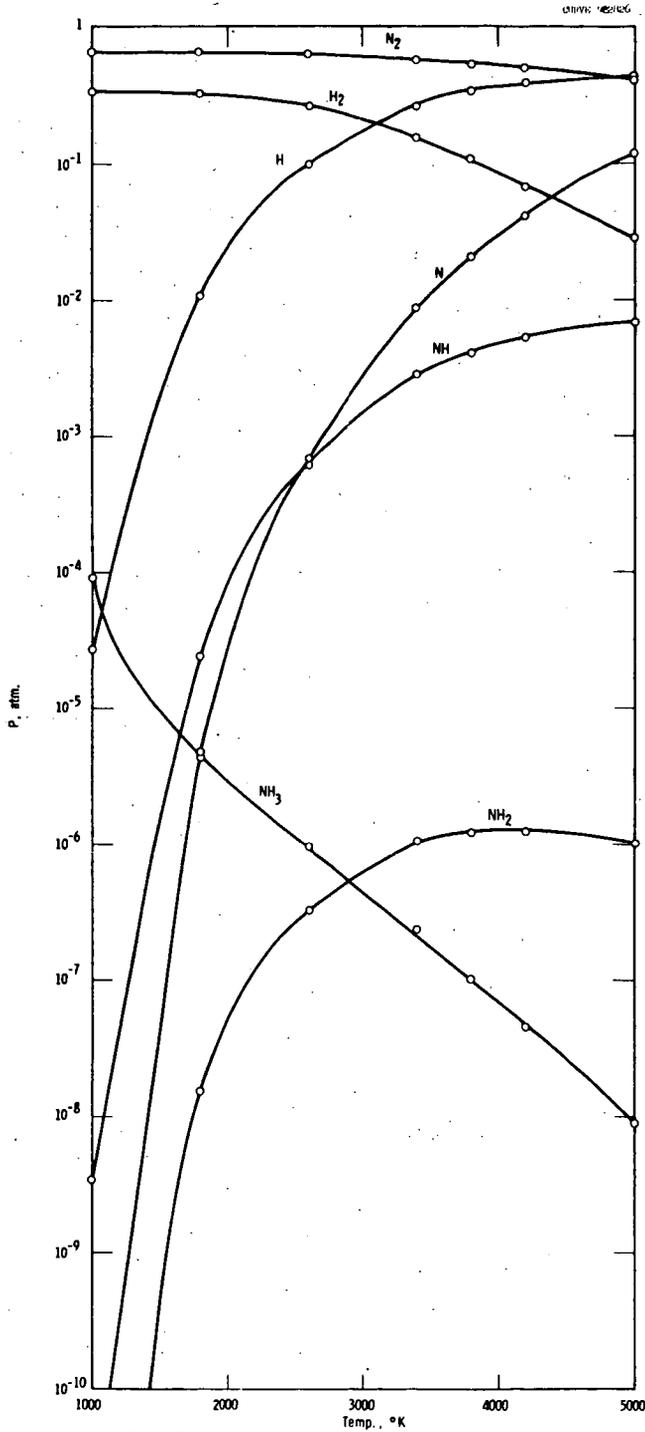


Fig. 1—Partial pressure as function of temperature for system  $2N_2 + H_2$  at total pressure of 1 atmosphere.

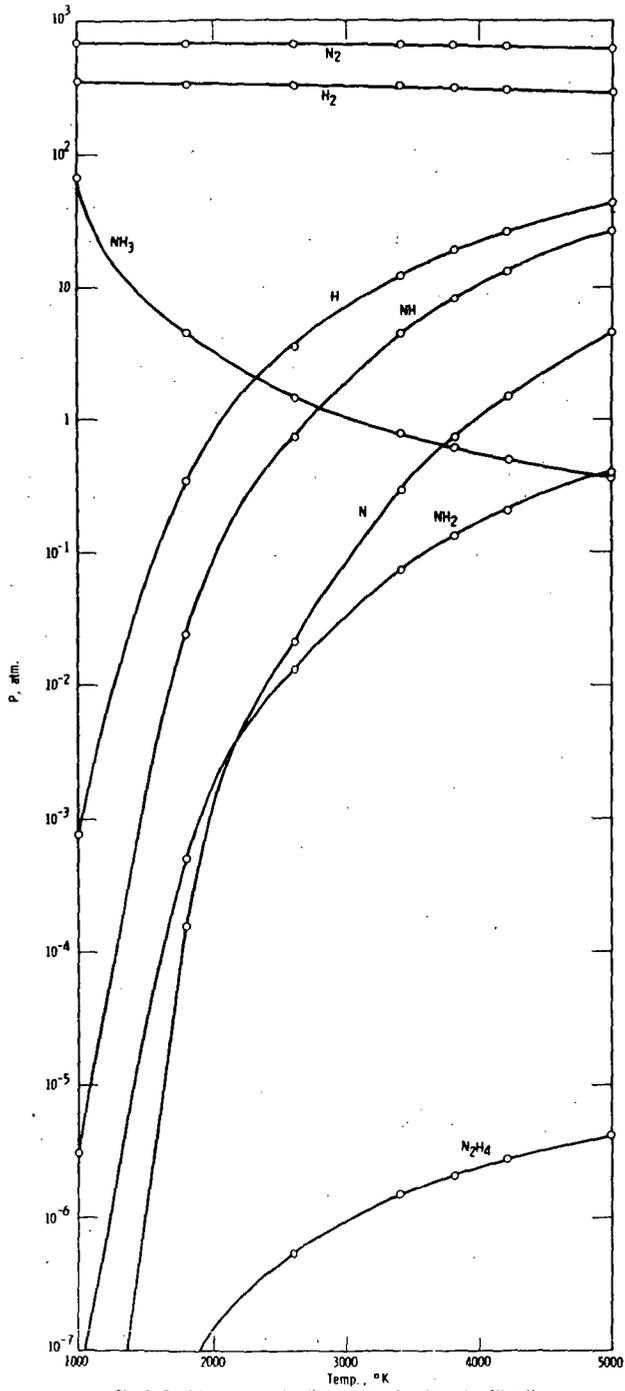


Fig. 2—Partial pressure as function of temperature for system  $2N_2 + H_2$  at total pressure of 1000 atmospheres.

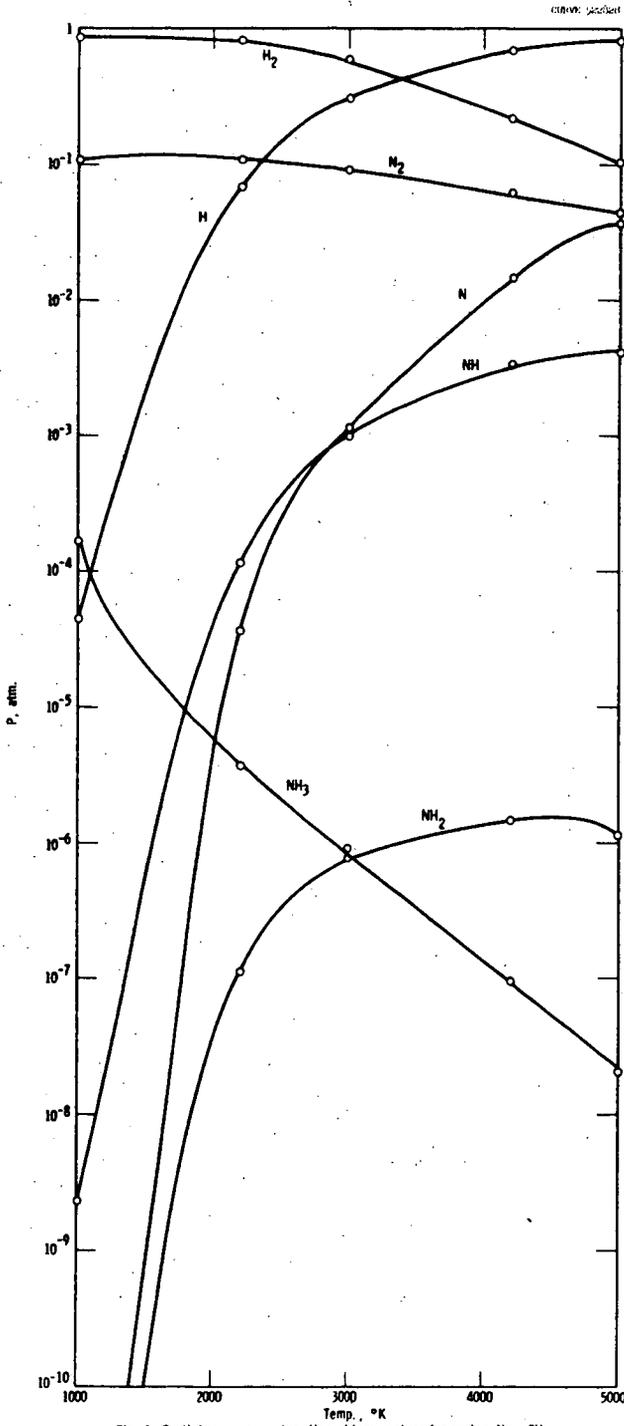


Fig. 3—Partial pressure as function of temperature for system  $N_2 + 3H_2$  at total pressure of 1 atmosphere.

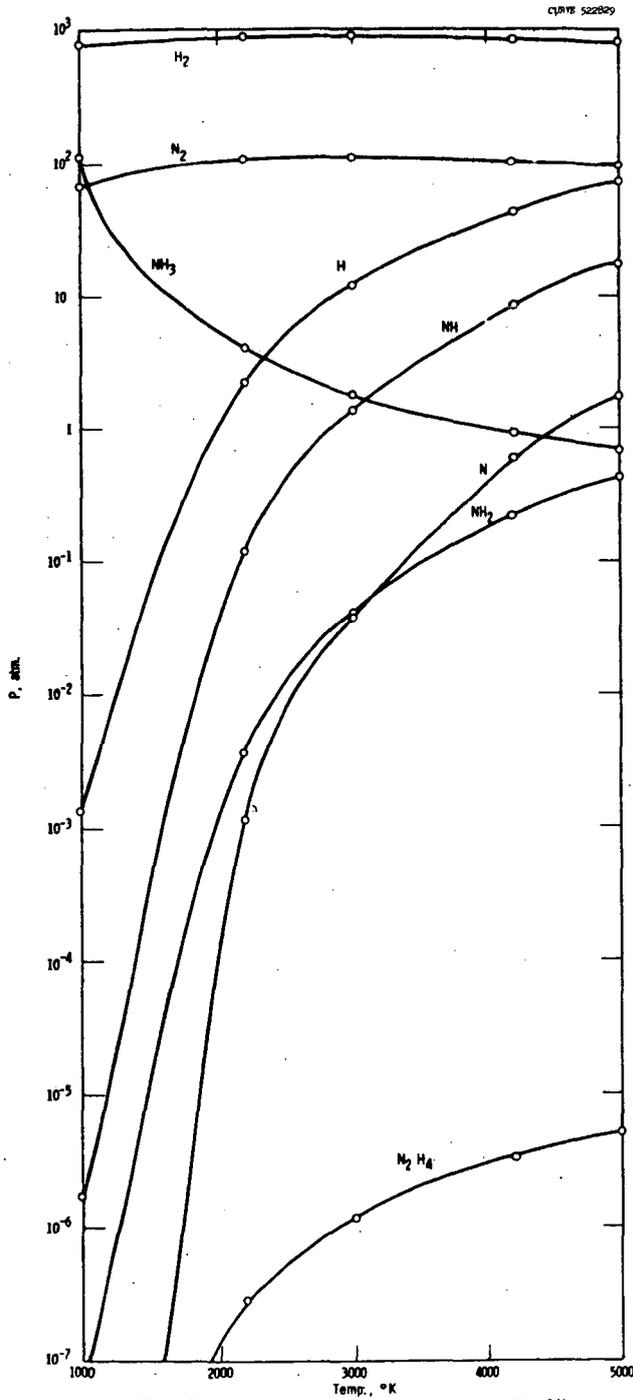


Fig. 4—Partial pressure as function of temperature for system  $N_2 + 8 H_2$  at total pressure of 1000 atmospheres.

TABLE I

MOLECULAR DATA FOR HYDRAZINE

<u>Vibrational Frequency</u>			
<u>Mode</u>	<u>cm<sup>-1</sup></u>	<u>Mode</u>	<u>cm<sup>-1</sup></u>
$\gamma_1$	3325	$\gamma_8$	3350
$\gamma_2$	3160	$\gamma_9$ (doublet)	3280 3314
$\gamma_3$	1493	$\gamma_{10}$	1628
$\gamma_4$	1098	$\gamma_{11}$	1275
$\gamma_5$	960	$\gamma_{12}$ (doublet)	933 966
$\gamma_6$	780		
$\gamma_7$	360		

Symmetry number ( $\sigma$ ) = 2

Molecular Weight = 32.048

Moments of Inertia

$$I_A = 6.18 \times 10^{-40} \text{ g cm}^2$$

$$I_B = 35.33 \times 10^{-40} \text{ g cm}^2$$

$$I_C = 36.98 \times 10^{-40} \text{ g cm}^2$$

Reduced moment for internal rotation =  $1.53 \times 10^{-40} \text{ g cm}^2$ 

Potential barrier height for internal rotation = 2800 cal/mole

$$\Delta H_0^{\circ} = 26,060 \text{ cal/mole}$$

Statistical weight of electronic ground state = 1

TABLE II

THERMODYNAMIC FUNCTIONS, FREE ENERGY OF FORMATION  
AND EQUILIBRIUM CONSTANT OF  $N_2H_4$  FROM 1000° TO 5000°K

Temp. °K	$C_p$ , Cal/Mole	$\frac{H^\circ - H_0^\circ}{T}$ k Cal/Mole Deg.	$S^\circ$ , Cal/Mole/Deg.	$\frac{F^\circ - H_0^\circ}{T}$ k Cal/Mole Deg.	$F_f^\circ$ , k Cal/Mole	$-\log K_p$
1000	21.95	15.21	76.06	61.64	77.18	16.89
1800	26.60	19.39	90.40	71.65	122.88	14.94
2600	28.50	21.92	100.54	79.24	167.83	14.12
3400	29.38	23.60	108.30	85.37	212.09	13.65
3800	29.64	24.21	111.58	88.01	234.14	13.48
4200	29.84	24.75	114.54	90.46	256.04	13.34
5000	30.12	25.57	119.73	94.36	299.54	13.11

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